

Design of a Selective Filter based on 2D Photonic Crystals Materials

Lallam Farah¹, Badaoui Hadjira², Abri Mehadji³

^{1,2}STIC Laboratory, Faculty of Technology, University of Tlemcen, Algeria

³Telecommunications Laboratory, Faculty of Technology, University of Tlemcen, Algeria

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ABSTRACT

Two dimensional finite differences temporal domain (2D-FDTD) numerical simulations are performed in cartesian coordinate system to determine the dispersion diagrams of transverse electric (TE) of a two-dimension photonic crystal (PC) with triangular lattice. The aim of this work is to design a filter with maximum spectral response close to the frequency 1.55 μm . To achieve this frequency, selective filters PC are formed by combination of three waveguides $W_1^k A$ wherein the air holes have of different normalized radii respectively $r_1/a=0.44$, $r_2/a=0.288$ and $r_3/a= 0.3292$ (a : is the periodicity of the lattice with value 0.48 μm). Best response is obtained when we insert three small cylindrical cavities (with normalized radius of 0.17) between the two half-planes of photonic crystal strong lateral confinement.

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Corresponding Author:

Lallam Farah,
Department of Physics,
Abu Bekr Belkaid University,
PB: 119 University of Tlemcen, 13000, Algeria.
Email: fa.lallam72@gmail.com

1. INTRODUCTION

Photonic crystals (PCs) are periodic dielectric nanostructures materials and have many applications: They also have applications in medical imaging field, solar cells and development of chemical and biological micro sensors, spectroscopy and electromagnetic shielding [1], [2]. The nanophotonic structures offer exceptional potential for the development of new architectures of photovoltaic solar cells in thin layers and high efficiency. Among the promising applications of PCs, we mentioned selective filter based on 2D CPs [3].

The PCs have a dielectric index which varies across the wavelength that is to be controlled, on one or more spatial directions. Analogous to electrons in semiconductors, the photon propagation can be described using a band structure in which transmission bands are separated by band gaps, energy ranges at which light cannot exist inside the photonic crystals [4]. Most of the research on PCs focuses on the use of band gaps. The insertion of defects is the easiest way to change the properties of photonic crystals and exploit them for the realization of amazing photonic components for integrated optics. Defects cause the appearance of modes within the band gap. Defects are realized by removal, addition or modification of pattern in one or more parallel rows of the crystal in ΓK direction of Brillouin first zone [5]. Hence, can be designed ($W_1^k A$) waveguide by creating between the two half-planes of photonic crystal strong lateral confinement, when single full row of air cylinders holes is removed in the PCs.

The aim of this work is to design a waveguide with maximum spectral response close to the frequency 1.55 μm . This nanostructure is composed by a superposition of three waveguides $W_1^k A$ with different radii. This kind of filter has been already studied in ref. [6]. The new in this paper is the insertion of

a number of defects between the two half-planes of photonic crystal strong lateral confinement, in order to improve the response of the filter.

In this paper, the simulation was based on the method of finite difference time domain (FDTD) [7]. 2-D numerical simulations are performed in Cartesian coordinate system to determine the dispersion diagrams of transverse electric (TE) of two-dimension PC with triangular lattice, operating as filter.

2. RESEARCH METHOD

The FDTD is used to numerically modeling a 2D photonic crystal containing air holes with triangular lattice. It is one of the most widely used numerical methods for computing the solution of electromagnetic problems, including photonic structures.

It provides us with a simple way to discretize the Maxwell's equations without requiring a complex mathematical formulation, and it does not require any symmetry in the structure being modelled, moreover FDTD can be used for the inhomogeneous structure materials in two or three dimension forms [7-10]. Furthermore, it computes the solution in the time domain, from which the frequency behaviour of the electromagnetic band gap elements can be extracted over a wide frequency range. The electromagnetic field is represented by Maxwell equations given by the following relations:

$$\begin{cases} \frac{\partial E_x}{\partial t} = \frac{1}{\epsilon} \left(\frac{\partial H_z}{\partial y} \right) \\ \frac{\partial E_y}{\partial t} = \frac{1}{\epsilon} \left(-\frac{\partial H_z}{\partial x} \right) \\ \frac{\partial H_z}{\partial t} = -\frac{1}{\mu} \left(\frac{\partial E_x}{\partial x} - \frac{\partial E_y}{\partial y} \right) \end{cases} \quad (1)$$

ϵ and μ represent respectively the permittivity and permeability of the material.

The FDTD method is used here to discretize these equations (1) with time step Δt , spatial steps Δx and Δy which are the distance between two neighboring grid points, respectively along the x and y directions in the xy-coordinate system. For the TE mode, the numerical scheme of the two-dimensional FDTD method [11] is:

$$\begin{aligned} \frac{E_x(x,y,t+\frac{\Delta t}{2}) - E_x(x,y,t-\frac{\Delta t}{2})}{\Delta t} &= \frac{1}{\epsilon} \left(\frac{H_z(x+\frac{\Delta y}{2},y,t) - H_z(x-\frac{\Delta y}{2},y,t)}{\Delta y} \right) \\ \frac{E_y(x,y,t+\frac{\Delta t}{2}) - E_y(x,y,t-\frac{\Delta t}{2})}{\Delta t} &= -\frac{1}{\epsilon} \left(\frac{H_z(x+\frac{\Delta x}{2},y,t) - H_z(x-\frac{\Delta x}{2},y,t)}{\Delta x} \right) \\ \frac{H_z(x,y,t+\frac{\Delta t}{2}) - H_z(x,y,t-\frac{\Delta t}{2})}{\Delta t} &= \frac{1}{\mu} \left(\frac{E_y(x+\frac{\Delta x}{2},y,t) - E_y(x-\frac{\Delta x}{2},y,t)}{\Delta x} - \frac{E_x(x,y+\frac{\Delta y}{2},t) - E_x(x,y-\frac{\Delta y}{2},t)}{\Delta y} \right) \end{aligned} \quad (2)$$

In order to achieve the numerical stability of the scheme [6], the time step value must satisfy the relation of the following criterion:

$$\Delta t \leq \frac{1}{c \sqrt{\left(\frac{1}{\Delta x}\right)^2 + \left(\frac{1}{\Delta y}\right)^2}} \quad (3)$$

Where c : is the velocity of light in free space. Also in our computation, the values of Δx and Δy respect the relation:

$$\Delta x = \Delta y \leq \frac{\lambda}{10\sqrt{\epsilon_r}} \quad (4)$$

Where ϵ_r : is the relative permittivity of the dielectric matrix and λ is the wavelength of the desired mode.

The computational domain has a rectangular shape in the x-y plane. The spatial discretization in the FDTD simulation is chosen to be $\Delta x = \Delta y = 0.04 \mu\text{m}$. A modulated Gaussian pulse is used to provide a wide-band excitation at any desired position inside the computational domain comprising the CP. Since the data storage in a computer is limited by the size of its memory, it is not possible to handle an open region problem directly. To mitigate this problem, the perfectly matched layer (PML) technique [12], [13] is widely used in

the FDTD simulations; it exhibits an accuracy level that is significantly better than most other absorbing boundary conditions (ABCs) [14], [15]. Here we consider only the conditions of absorption-type wall that simulate a finite domain containing the entire structure study by investigating the lowest reflection digital interfaces.

The adopted two-dimensional finite difference time domain simulation process is shown in Figure 1 in xyz coordinate system. According to the Figure 1, two located detectors are used to compute the reflection at the input and transmission at the output of the structure. A Gaussian source is employed for excitation (source). The structure is surrounded by four surrounding absorbing walls.

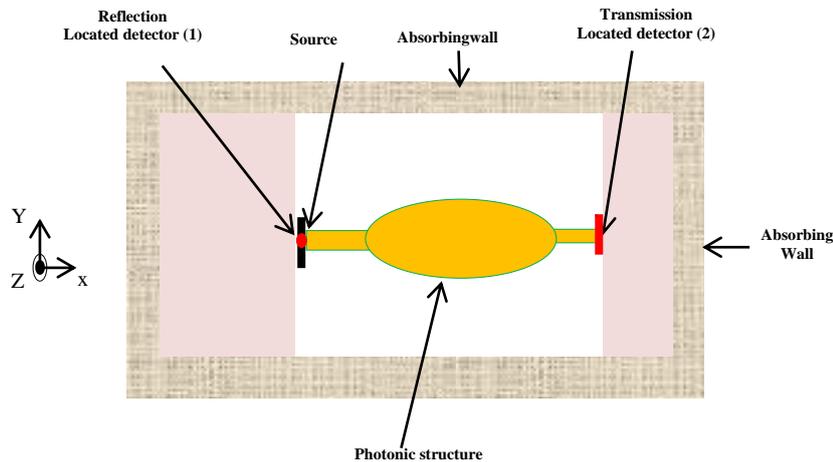


Figure 1. Two dimensional finite difference time domain simulation process

3. RESULTS AND ANALYSIS

In the following, numerical simulation is done for the first topology which is the combination of three kinds of waveguides $W_1^K A$ coupled in cascade arrangement within the same cell of PC with triangular lattice and without cavities inside the two half-planes of photonic crystal strong lateral confinement (Figure2).

Each waveguide possesses a number of holes $n=15$ (total number of holes = 45). The normalized size of holes are respectively $r_1/a = 0.44$, $r_2/a = 0.288$, $r_3/a = 0.3292$ and the lattice constant $a = 0.48\mu\text{m}$. The air cylinders holes having are embedded in a dielectric matrix doped INP/GAINASP/INP; refractive index of matrix is $n = 3.24$. These air rods are arranged in the x direction and supposed infinitely long in the z direction.

The transmission coefficient versus wavelength for the TE polarization, derived from the FDTD simulation corresponding to the first topology is plotted in Figure 3(a). We observe that the response of the filter is peaked around $1.55\mu\text{m}$. The maximal quantity of transmission computed is 77%. Others peaks appear very close of this frequency with coefficient of transmission of value 80% ranged in the frequency band $[1.54-1.57]\mu\text{m}$ (Figure 3(b)).

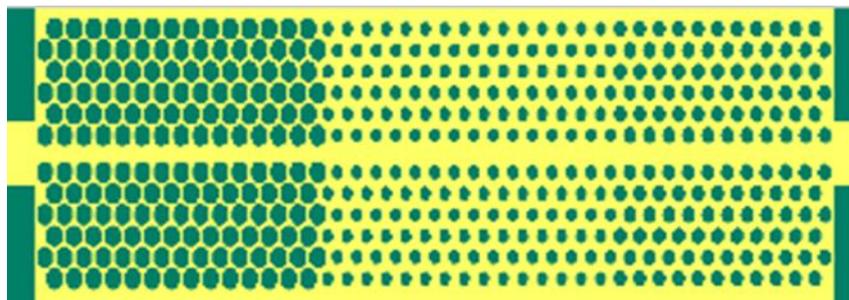


Figure 2. Modeling scheme of PC with combination of three guides $W_1^K A$ with triangular lattice and without cavities

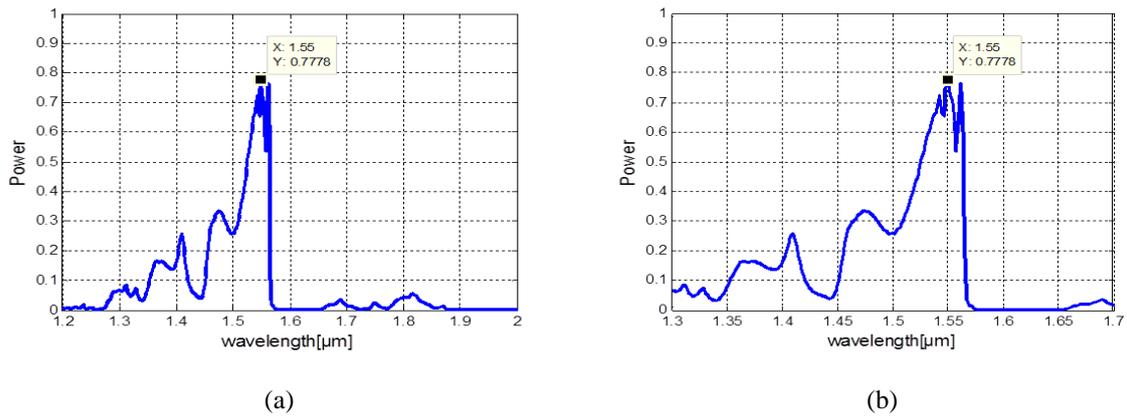


Figure 3. Computed transmission coefficients for the three simulated selective filter (without cavities). $r_1/a=0.44$, $r_2/a=0.288$ and $r_3/a=0.3292$ (a : is the periodic lattice of value $0.48 \mu\text{m}$)

Another two-dimensional PC filter is simulated in order to eliminate unwanted peaks and filter the desired frequency, based on the same first topology with the three $W_1^K A$ but we add three small cylindrical holes with normalized radius of 0.17 between the two half-planes of photonic crystal strong lateral confinement. All the previous geometrical parameters are unchanged (Figure 4). We report in the Figure 5 the transmission coefficient versus wavelength for the TE polarization. According to this plot, it can be seen apparently that maximum spectral response (80%) occurs close to the frequency $1.55 \mu\text{m}$. A significant improvement was observed regarding the disappearance of unwanted peaks around the desired frequency but a slight decrease occurs in the amplitude of appearing modes.

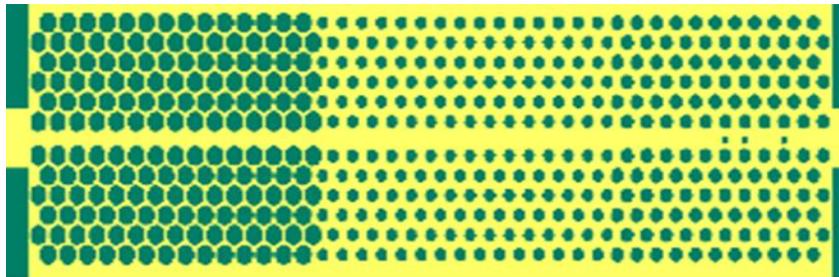


Figure 4. Modelling scheme of PC with combination of three guides $W_1^K A$ with triangular lattice and having three cavities

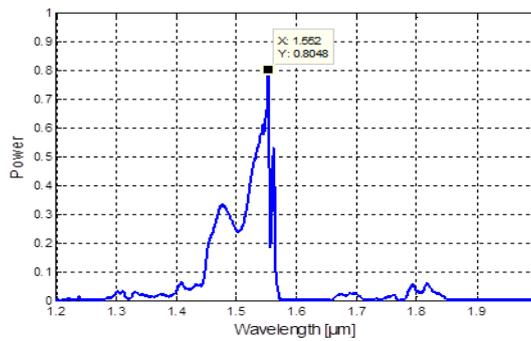
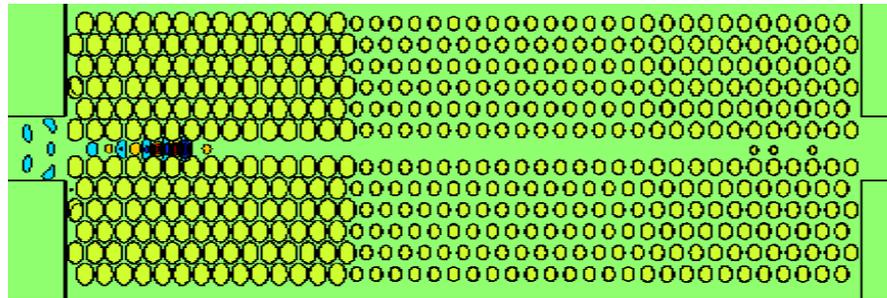
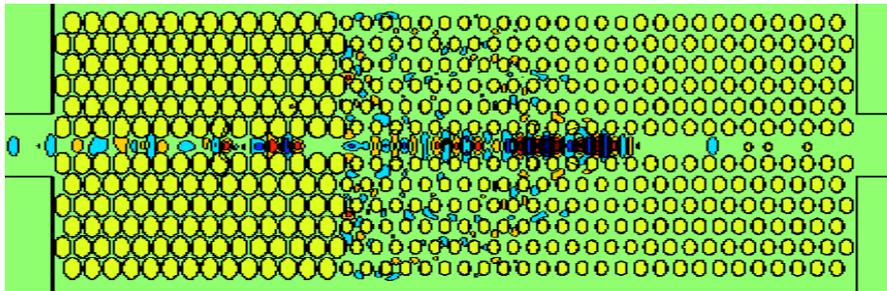


Figure 5. Computed transmission coefficients for the three simulated selective filter with three cavities having normalized radius of 0.17. $r_1/a=0.44$, $r_2/a=0.288$ and $r_3/a=0.3292$ (a is the periodic lattice of value $0.48 \mu\text{m}$)

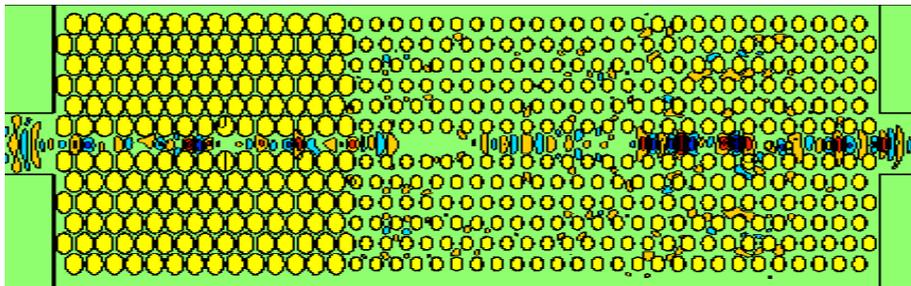
The magnetic field pattern inside the selective filter corresponding to the second topology with a spectral response close to frequency of $1.55\mu\text{m}$ is reported in Figure 6(a-c) for different step time iterations 1500, 4500 and 6500. These figures show the light-guiding propagation of the electromagnetic field inside the empty row along the waveguide (confinement). We can observe that one part of electromagnetic energy is transmitted until the end of PC with a frequency belonging to the gap, and the other part of energy with a non-allowed frequency is reflected by the added inclusions embedded in the dielectric matrix in the empty row.



(a)



(b)



(c)

Figure 6. Simulated distribution of a magnetic field. (a) 1500, (b) 4500, (c) 6500 step time iterations

4. CONCLUSION

A novel selective filter based on a two-dimensional photonic crystal with a triangular lattice with air holes in a dielectric substrate was proposed. This device consists of three waveguides W_1^K placed in cascade with different radii respectively $r_1/a=0.44$, $r_2/a=0.288$ and $r_3/a=0.3292$. An improvement is observed when we add a number of defects in the empty row. From our computation, this CP can be conceived for achieving a frequency of $1.55\mu\text{m}$ which is suitable for optical devices with a transmission maximum in order to 80%.

In order to improve more the selectivity of this filter, other simulations with new designs will be done later. This study will be completed when comparison with experimental data is done.

REFERENCES

- [1] L. Jylhä, I. Kolmakov, S. Maslovski, S. Tretyakov, "Modeling of Isotropic Backward -wave Materials Composed of Resonant Spheres," *Journal of Applied Physics*, pp 99, 043102, 2006.
- [2] F. D. Mahad, Abu Sahmah M. Supa'at, D. Forsyth, T. Sun, A. Izam Azmi, "Characterization of Erbium Doped Photonic Crystal Fiber," *TELKOMNIKA*, vol. 14, no.3, pp. 880-886, 2016.
- [3] R. D. Meade, A. Devenyi, J. D. Joannopoulos, O. L. Alerhand, D. A. Smith et K. Kash, "Novel Applications of Photonic band Gap Materials: Low Loss bends and Q Cavities," *Journal of Applied Physics*, 75, pp. 4753, 1994.
- [4] E. Yablonovitch, "Inhibited Spontaneous Emission in Solid-State Physics and Electronics," *Phys. Rev. Lett.*, 58, pp. 2059-2062, 1987.
- [5] J. D. Joannopoulos, R. D. Maede, J. N. Winn, "Photonic crystals: Modelling the row of light," *Princeton Univ. Press*, 1995.
- [6] H. Badaoui and M. Abri, "New Design of Integrated 2D Photonic Crystal Narrow Band Filters Using the FDTD-2D Method," *Frequenz*, 68(11-12): pp. 511-518, 2014.
- [7] A. Taflove, "Advances in Computational Electrodynamics: The Finite-Difference Time-Domain Method," Norwood, MA: Artech House, 1998.
- [8] Hao, Y., C. J. Railton, "Analyzing Electromagnetic Structures with Curved Boundaries on Cartesian FDTD Meshes," *IEEE Trans. Microwave Theory tech.* 46, pp. 82-88, 1998.
- [9] Luebbers, R. et al., "A Frequency-Dependent Finite-Difference Time-Domain Formulation for Dispersive Materials," *IEEE Trans. Electromagn. Compat.* 32, pp. 222-227, 1990.
- [10] Baohe Yuan, Qi Xu, Linfei Liu, Xiyang Ma, "Investigation Effects of Filling Rate on the Bands Gaps of Two Dimensional Photonic Crystals," *TELKOMNIKA*, vol. 14, no. 3A, pp. 27-32, 2016.
- [11] V. Liu, Y. J.P. Berenger, "A Perfect Matched Layer for the Absorption of Electromagnetic waves," *J. Comput. Phys.*, 114, pp. 185-200 1994
- [12] Jiao, David A. B. Miller, S. Fan, "Design Methodology for Compact Photonic-Crystal-based Wavelength Division Multiplexers," *Optics Letters*, 36 (4), 2011.
- [13] J. C. Veihl, R. Mittra, "An Efficient Implementation of Berengers Perfectly Matched Layer (PML) for Finite-difference Time-domain mesh Truncation," *IEEE Microwave Guided Wave Lett.*, 6, pp. 9496, 1996.
- [14] B. Engquist, A. Majda, "Absorbing Boundary Conditions for the Numerical Simulations of Waves," *Math. Comp.*, 31, pp. 629-651 1977.
- [15] G. Mur, "Absorbing Boundary Conditions for the Finite difference Approximation of the Time-domain Electromagnetic-field Equations," *IEEE Trans. Electromagn. Compat.*, 23, pp. 377-382 1981.